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## TRIPLET STATES OF OLIGOTHIOPHENES AND ELECTRON TRANSFER TO C<sub>60</sub>: A PULSED EPR INVESTIGATION

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**Abstract** Triplet states of oligothiophenes and C<sub>60</sub> were excited by a pulsed laser light using a Nd:YAG laser. The finestructure values D, E of different endcapped oligothiophenes (EC3T, EC4T, and EC5T) were determined from pulsed EPR-spectra. Time resolution was applied to study the kinetics and relaxation times.

A photoinduced electron transfer from EC4T to C<sub>60</sub> in liquid solutions was investigated in detail. Fourier transform of the free induction decays results in spectra of the transient triplet state and of the transient C<sub>60</sub> monoradical anion. Electron transfer rates were deduced from the observed triplet quenching rates at different temperatures.

## INTRODUCTION

The investigation of the photoexcitation in oligomers molecules with well defined chain length represents a useful approach for the understanding of the electronic properties in the corresponding polymers. Particularly time resolved EPR spectroscopy can give informations about the nature, the relaxation processes and the lifetime of an electronic excitation.

The end-capped thiophene oligomers<sup>1</sup> are of high interest as model compounds for polythiophene and for photoinduced electron transfer processes with electron acceptor molecules<sup>2</sup>. Spectroelectrochemical investigations show that oxidation-potentials and the  $\pi - \pi^*$  excitation energy vary systematically with the chain length. Furthermore the compounds with chain length  $n \geq 4$  undergo two stable, reversible oxidized states at room temperature<sup>3</sup>. C<sub>60</sub> fullerene on the other hand represents also a new molecular system with very unusual electronic properties. After light excitation the lowest triplet state of the molecule is populated with a quantum yield of nearly one which indicates that the system would be an ideal component in photochemical devices. The electron affinity is so high that C<sub>60</sub> can be reduced chemically as well as

electrochemically with up to six electrons<sup>4</sup>. In order to study the properties of some end-capped oligothiophenes ( $n = 3 - 5$ ) and  $C_{60}$  as a new donor/acceptor system we performed a pulsed EPR investigation of their lowest photoexcited triplet states in glassy matrix. Furthermore we investigated the dynamics of an electron transfer process from EC4T to  $^3C_{60}$  in liquid solution.

## EXPERIMENTAL

Toluene solutions ( $c \approx 1 \cdot 10^{-3}$  mol/l) of the end-capped oligothiophenes (EC3T, EC4T, EC5T),  $C_{60}$  and  $C_{60}$  in mixture with EC4T were degassed by several freeze-thaw cycles and sealed off on a high vacuum line. The samples were irradiated at  $\lambda = 355$  nm by a Nd:YAG-laser with an energy of  $\simeq 20$  mJ per pulse at a pulse frequency of 30 Hz.

The EPR measurements were carried out using a home built electron spin echo spectrometer with a slotted tube resonator and an helium flow cryostat. The triplet EPR powder spectra were recorded at  $T = 10$  K with a Hahn-echo microwave pulse sequence by integrating the echo signal while sweeping the magnetic field. The lineshape of the spectra corresponds directly to the EPR line and should not be confused with standard EPR spectra recorded with field modulation. The time evolution of the EPR-line was obtained by varying the delay time between laser and first microwave pulse. Free induction decays were recorded using quadrature detection. The short deadtime of the spectrometer after the microwave pulse ( $\tau_D \simeq 70$  ns) allowed to Fourier transform the data set to give the EPR spectrum after a frequency proportional phase correction.

## OLIGOTHIOPHENE TRIPLETS

The samples containing the glassy toluene solutions of the end-capped oligomers were irradiated at  $T = 10$  K. We observe highly polarized EPR spectra in all samples when starting the Hahn-echo detection sequence at  $t \simeq 20$  ns after the laser pulse. The spectra are represented in figure 1. All signals grow up in a time scale shorter than the spectrometer time resolution. This is expected because the lowest excited triplet state of the molecule is populated during the laser flash via highly excited singlet states by mean of the spin orbit interaction in a time scale of  $\tau_{ISC} = 10^{-8} - 10^{-9}$  s. The intersystem crossing causes a selective population of the zero field energy levels which depends on the molecular symmetry. Since the high field Zeeman wave functions are linear combinations of the zero field wave functions the population of the triplet Zeeman levels do not correspond to the Boltzmann equilibrium population but shows polarization patterns typical for the given molecular symmetry.

In this case all spectra show the same polarization pattern, i.e. emissive peaks at the low field side and absorptive peaks at the high field side. The polarization was determined by comparing the triplet signal phase with the one of a stable radical measured simultaneously.

Triplet EPR powder spectra are characterized by their total spectral width which is correlated with the finestructure parameter  $D$ . We notice that the spectral width (edge to edge) decreases with increasing oligomers chain lengths. This indicates that we are observing the intrinsic triplet state on the molecule and not a local excitation of a thiophene ring. The finestructure parameters  $D$  and  $E$  which characterize the strength and the asymmetry of the electron dipolar coupling can not be determined directly from the spectra because the high field approximation is not valid for large  $D$  and  $E$  values. Therefore we simulated the EPR powder spectra using the solution for the eigenvalues of the Hamilton-operator:

$$\hat{H} = g\beta\mathbf{SB} - (XS_x^2 + YS_y^2 + ZS_z^2)$$

where  $X$ ,  $Y$ ,  $Z$  are the eigenvalues of the finestructure tensor in the tensor principal axis. The polarization for every transition was calculated from the population of the zero field energy levels, the transition probabilities were considered for a microwave field perpendicular to the magnetic field. The lineshape of the powder spectra computed for a static triplet do not correspond exactly to the experimental one. Obviously the intensity of transitions in orientations between the principal axis is much weaker than expected for a powder pattern. Measurements of the spectra at different delay time after the laser pulse show that the low field side decays much faster than the high field side and the spectra become absorptive after a time delay of  $\tau \simeq 30 \mu\text{s}$ . This indicates that the spin lattice relaxation is on the order of few microseconds and may also affect the spectra recorded 500 ns after the laser pulse. For this reason we considered anisotropic  $T_1$  relaxation in the spectra calculation. Spin lattice relaxation in a three level system is biexponential with the rates  $k_1 = 3w$  and  $k_2 = -(w+2W)$ .  $w$  and  $W$  are the probabilities for  $\Delta m = 1$  and  $\Delta m = 2$  transitions respectively. The angular dependence of  $w$  and  $W$  can be obtained from the matrix elements of the Hamilton-operator in the laboratory axis system which are proportional to the  $S_{\pm}$  and  $S_{\pm}^2$  operators respectively. For the simulations at low temperature it can be assumed that  $\omega_0\tau_c \gg 1$  and therefore the spectral densities for  $w$  and  $W$  are related by  $J_1 \simeq 4J_2$ . The spectra simulations are represented in figure 1 by the solid lines. The obtained finestructure values and zero field populations are listed in table 1.

TABLE 1: Triplet state finestructure values and populations of EC3T, EC4T, EC5T

	$ D  [cm^{-1}]$	$ E  [cm^{-1}]$	$P_x$	$P_y$	$P_z$
EC3T	0.0738	0.0046	0.31	0.18	0.51
EC4T	0.0682	0.0023	0.35	0.21	0.44
EC5T	0.0607	0.0023	0.35	0.25	0.4

It is not possible to determine the absolute sign of the  $D$  value from the calculation. Considering the particular structure of the oligomers with a stretched one-dimensional  $\pi$ -system we suggest a prolate triplet spin density distribution which corresponds to a negative  $D$ -value.

Comparing the experimental with the calculated spectra we notice small deviations in the case of the EC3T. This could be due to the anisotropic hyperfine interaction which can lead to anisotropy in the linewidth. A computation of the spectra including hyperfine coupling may be performed if the anisotropic hyperfine tensor (from single crystal measurements) is known which is currently not the case.

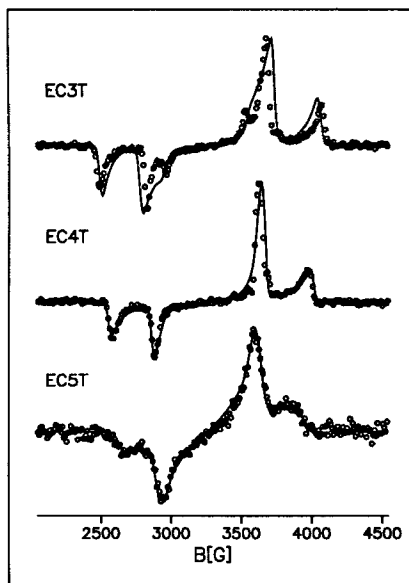


FIGURE 1: Triplet EPR powder spectra of the end-capped oligothiophenes EC3T, EC4T, EC5T

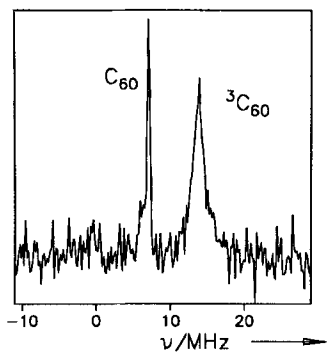


FIGURE 2: Fourier transform spectrum of  ${}^3\text{C}_{60}$  and  $\text{C}_{60}^-$

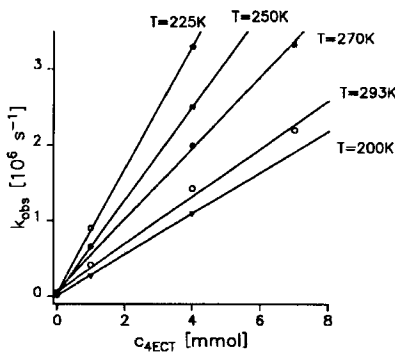


FIGURE 3: Stern-Volmer plots of the observed triplet decay rates

### ELECTRON TRANSFER FROM EC4T TO ${}^3\text{C}_{60}$

Toluene solutions containing mixtures of EC4T and  $\text{C}_{60}$  in different concentration ratios were irradiated in a temperature range  $T = 200 - 300$  K. A free induction decay (FID) was observed after the laser pulse. The Fourier transform of the FID gives a spectrum with two narrow absorptive EPR lines at  $g = 2.0014$  and  $g = 1.9999$ . The EPR spectrum is shown in figure 2. We assign the first line to the triplet state of  $\text{C}_{60}$  and the second line to  $\text{C}_{60}$  radical monoanion according to the following analysis of the kinetics.

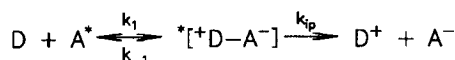
The build-up and decay of the  $\text{C}_{60}$  triplet EPR line can be described by a biexponential relation which considers only  $T_1$  relaxation and pseudo first order kinetics for the electron transfer as dominant processes. The time evolution of the triplet EPR line at different temperatures was fitted with the equation:

$$S(t) \sim M_B(-\exp[-t/T_1] + \exp[-k_{obs}t])$$

We observe that the triplet decay rate  $k_{obs}$  depends on the concentration of the added donor EC4T. When plotting  $k_{obs}$  versus the concentration of EC4T we obtain for a given temperature a linear correlation (figure 3). The quenching follows the Stern- Volmer relation:

$$k_{obs} = k_T + k_q[\text{EC4T}]$$

where  $k_T$  is the triplet decay rate without quencher. The determined quenching rate  $k_q$  increases with decreasing temperature. We estimated a negative activation energy of  $E_a \simeq -10$  kJ/mol which is in contrast to the model of a diffusion controlled or thermally activated electron transfer process according to the theory of Marcus. A second quenching mechanism has to be considered which includes the formation of an exciplex intermediate before the charge separation takes place<sup>5</sup>:



In this path the equilibrium constant  $K_{exc} = k_1/k_{-1}$  involves exciplex dissociation which for a bound complex is expected to be strongly temperature dependent. If the exciplex formation enthalpy is larger than the ion pair solvation enthalpy ( $\Delta H_{ip} < -\Delta H_{exc}$ ) the activation energy is indeed negative as observed here.

The second component of the EPR spectrum has a g-factor which corresponds to that found in electrochemically prepared  $C_{60}$  radical monoanions<sup>4</sup>. The assignment of this transient to  $C_{60}^-$  is also supported by the analysis of the signal polarization. The signal is absorptive at all times as is expected from the radical pair mechanism: the polarization of radicals formed from a non polarized triplet exciplex is positive for the transient with the lower g-factor<sup>6</sup>, i.e. for  $C_{60}^-$  (compared to  $g = 2.0024$  for EC4T). We determined the temperature dependence of the linewidth for  $C_{60}^-$  and  ${}^3C_{60}$  by fitting the EPR lines with a Lorentz line shape function. The  $1/T_2$  values are plotted versus temperature in figure 4.

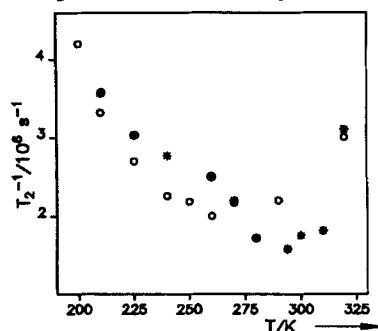


FIGURE 4: Temperature dependence of the EPR-linewidth ( $1/T_2$ ) of (o)  ${}^3C_{60}$  and (\*)  $C_{60}^-$

The relaxation rates of the two transients show the same temperature dependence. This suggests that  $T_1$ ,  $T_2$  are determined by the same relaxation mechanism for both  ${}^3C_{60}$  and  $C_{60}^-$  which we attribute to the spin rotation interaction.

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